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The integrals in Gradshteyn and Ryzhik. Part 20: Hypergeometric functions

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ABSTRACT. The table of Gradshteyn and Ryzhik contains many integrals that involve the hypergeometric function ${}_pF_q$. Some examples are discussed.

1. Introduction

The hypergeometric function defined by

$$(1.1) \quad {}_pF_q(a_1, a_2, \dots, a_p; b_1, b_2, \dots, b_q; x) := \sum_{k=0}^{\infty} \frac{(a_1)_k \cdots (a_p)_k}{(b_1)_k \cdots (b_q)_k} \frac{x^k}{k!}$$

includes, as special cases, many of the elementary special functions. For example,

$$(1.2) \quad \begin{aligned} \log(1+x) &= x {}_2F_1(1, 1; 2; -x) \\ \sin x &= x {}_0F_1\left(-; \frac{3}{2}; -x^2/4\right) \\ \cosh x &= \lim_{a, b \rightarrow \infty} {}_2F_1\left(a, b; \frac{1}{2}; x^2/4ab\right). \end{aligned}$$

The binomial theorem, for real exponent, can also be expressed in hypergeometric form as

$$(1.3) \quad (1-x)^{-a} = {}_1F_0(a; -; x).$$

The goal of this paper is to verify the integrals in [3] that involve this function. Due to the large number of entries in [3] that can be related to hypergeometric functions, the list presented here represents the first part of these. More entries will appear in a future publication.

The hypergeometric function satisfies a large number of identities. The reader will find in [1] the best introduction to the subject. Some elementary identities are

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described here in detail. For example, if one of the top parameters (the a_i) agrees with a bottom one (the b_i), the function reduces to one with lower indices. The identity

$$(1.4) \quad {}_2F_1(a, b; a; x) = {}_1F_0(a; -; x).$$

illustrates this point. The binomial theorem identifies the latter as $(1-x)^{-a}$.

2. Integrals over $[0, 1]$

The first result is a representation of ${}_2F_1$ in terms of the *beta integral*

$$(2.1) \quad B(a, b) = \int_0^1 t^{a-1}(1-t)^{b-1} dt.$$

Proposition 2.1. The hypergeometric function ${}_2F_1$ is given by

$$(2.2) \quad {}_2F_1(a, b; c; x) = \frac{1}{B(b, c-b)} \int_0^1 t^{b-1}(1-t)^{c-b-1}(1-tx)^{-a} dt.$$

PROOF. Expand the term $(1-tx)^{-a}$ by the binomial theorem and integrate term by term. \square

This representation appears as **3.197.3** in [3]. In order to simplify the replacing of parameters, this entry is also written as

$$(2.3) \quad \int_0^1 t^b(1-t)^c(1-tx)^a dt = B(b+1, c+1) {}_2F_1(-a, b+1; b+c+2; x).$$

This is one of the forms in which it will be used here: the integral being the object of primary interest.

Example 2.2. The special case $a = c = 1$ in (2.2) appears as **3.197.10** in [3]:

$$(2.4) \quad \int_0^1 \frac{t^{b-1} dt}{(1-t)^b(1+tx)} = \frac{\pi}{\sin \pi b} (1+x)^{-b}.$$

The evaluation is direct. The identity (1.4) gives

$$(2.5) \quad {}_2F_1(1, b; 1; -x) = (1+x)^{-b}$$

and then use $B(b, 1-b) = \Gamma(b)\Gamma(1-b) = \pi/\sin \pi b$ to complete the evaluation.

Example 2.3. Introduce the index r by $r = a - b$ and take $c = b + r$ in (2.2). Then we have

$$(2.6) \quad \int_0^1 t^{b-1}(1-t)^{r-1}(1-tx)^{-b-r} dt = B(b, r) {}_2F_1(b+r, b; b+r; x)$$

The identity (1.4) reduces the previous evaluation to

$$(2.7) \quad \int_0^1 t^{b-1}(1-t)^{r-1}(1-tx)^{-b-r} dt = B(b, r) (1-x)^{-b}.$$

This appears as **3.197.4** in [3].

3. A linear scaling

In this section integrals obtained from the basic representation (2.3) by the change of variables $y = tp$. This produces

$$(3.1) \quad \int_0^p y^{b-1}(p-y)^{c-b-1}(p-xy)^{-a} dy = p^{c-a-1} B(b, c-b) {}_2F_1(a, b; c; x).$$

Example 3.1. The special case $c = b + 1$ produces

$$(3.2) \quad \int_0^p y^{b-1}(p-xy)^{-a} dy = \frac{1}{b} p^{b-a} {}_2F_1(a, b; b+1; x),$$

where we have used $B(b, 1) = 1/b$. In order to eliminate the factor p^{-a} , we choose $x = -pr$ to obtain

$$(3.3) \quad \int_0^p y^{b-1}(1+ry)^{-a} dy = \frac{1}{p} p^b {}_2F_1(a, b; b+1; -rp),$$

This appears as **3.194.1** in [3]. The special case $a = 1$, stating that

$$(3.4) \quad \int_0^p \frac{y^{b-1} dy}{1+ry} = \frac{1}{b} p^b {}_2F_1(1, b; b+1; -rp),$$

appears as **3.194.5** in [3].

Example 3.2. The table [3] contains the formula **3.196.1**:

$$(3.5) \quad \int_0^u (x+b)^\nu (u-x)^{\mu-1} dx = \frac{b^\nu u^\mu}{\mu} {}_2F_1\left[1, -\nu, 1+\mu, -\frac{u}{b}\right].$$

We believe that it is a bad idea to have u and μ in the same formula, so we write this as

$$(3.6) \quad \int_0^a (x+b)^\nu (a-x)^{\mu-1} dx = \frac{b^\nu a^\mu}{\mu} {}_2F_1\left[1, -\nu, 1+\mu, -\frac{a}{b}\right].$$

To prove this, we let $x = at$ to get

$$(3.7) \quad \int_0^a (x+b)^\nu (a-x)^{\mu-1} dx = b^\nu a^\mu \int_0^1 (1+at/b)^\nu (1-t)^{\mu-1} dt.$$

The integral representation (2.3) now gives the result.

4. Powers of linear factors

The hypergeometric function appears in the evaluation of integrals of the form

$$(4.1) \quad I = \int_a^b L_1(x)^{\mu-1} L_2(x)^{\nu-1} L_3(x)^{\lambda-1} dx$$

where L_j are linear functions and $L_1(a) = L_2(b) = 0$. For example, **3.198**:

$$(4.2) \quad \int_0^1 x^{\mu-1} (1-x)^{\nu-1} [ax + b(1-x) + c]^{-(\mu+\nu)} dx = (a+c)^{-\mu} (b+c)^{-\nu} B(\mu, \nu)$$

is reduced to the normal form (2.3) by writing

$$(4.3) \quad I = (b+c)^{-\mu-\nu} \int_0^1 x^{\mu-1} (1-x)^{\nu-1} (1-rx)^{-(\mu+\nu)} dx$$

with $r = (b-a)/(b+c)$. Then (2.3) gives

$$(4.4) \quad I = (b+c)^{-\mu-\nu} B(\mu, \nu) {}_2F_1 \left(\mu + \nu, \mu; \mu + \nu; \frac{b-a}{b+c} \right).$$

To produce the stated answer, simply observe the special value of the hypergeometric function

$$(4.5) \quad {}_2F_1(a, b; a; z) = (1-z)^{-b}.$$

Similarly, the evaluation of **3.199**:

$$(4.6) \quad \int_a^b (x-a)^{\mu-1} (b-x)^{\nu-1} (x-c)^{-\mu-\nu} dx = (b-a)^{\mu+\nu-1} (b-c)^{-\mu} (a-c)^{-\nu} B(\mu, \nu),$$

is reduced to the interval $[0, 1]$ by $t = (x-a)/(b-a)$ and then the result follows from **3.198**.

The specific form of the answer is sometimes simplified due to a special relation of the parameters μ , ν and λ in (4.1). For example, in the evaluation of **3.197.11**:

$$(4.7) \quad \int_0^1 \frac{x^{p-1/2} dx}{(1-x)^p (1+qx)^p} = \frac{2}{\sqrt{\pi}} \Gamma(p + \frac{1}{2}) \Gamma(1-p) \cos^{2p}(\varphi) \frac{\sin((2p-1)\varphi)}{(2p-1) \sin(\varphi)},$$

with $\varphi = \arctan \sqrt{q}$. The standard reduction of the integral to hypergeometric form is easy. Write

$$(4.8) \quad I = \int_0^1 x^{p-1/2} (1-x)^{-p} (1+qx)^{-p} dx$$

and use (2.3) to obtain

$$(4.9) \quad I = B(p + \frac{1}{2}, 1-p) {}_2F_1(p, p + \frac{1}{2}; \frac{3}{2}; -q).$$

To reduce the answer to the stated form, we employ **9.121.19**:

$${}_2F_1 \left(\frac{n+2}{2}, \frac{n+1}{2}; \frac{3}{2}; -\tan^2 z \right) = \frac{\sin nz \cos^{n+1} z}{n \sin z}.$$

The evaluation of **3.197.12**:

$$(4.10) \quad \int_0^1 \frac{x^{p-1/2} dx}{(1-x)^p (1-qx)^p} = \frac{\Gamma(p + \frac{1}{2}) \Gamma(1-p)}{\sqrt{\pi}} \frac{[(1-\sqrt{q})^{1-2p} - (1+2\sqrt{q})^{1-2q}]}{(2p-1) \sqrt{q}}.$$

is done in similar form. The reduction to

$$(4.11) \quad I = B(p + \frac{1}{2}, 1-p) {}_2F_1(p, p + \frac{1}{2}; \frac{3}{2}; q)$$

is direct from (2.3). The stated form now follows from **9.121.4**:

$${}_2F_1 \left(-\frac{n-1}{2}, -\frac{n}{2} + 1; \frac{3}{2}; \frac{z^2}{t^2} \right) = \frac{(t+z)^n - (t-z)^n}{2nzt^{n-1}}.$$

5. Some quadratic factors

The table [3] contains several entries of the form

$$(5.1) \quad I = \int_a^b Q_1(x)^{\mu-1} L_2(x)^{\nu-1} L_3(x)^{\lambda-1} dx$$

where $Q_1(x)$ is a quadratic polynomial and L_j are linear functions. These are discussed in this section.

Example 5.1. The first entry evaluated here is **3.254.1**

$$\int_0^a x^{\lambda-1} (a-x)^{\mu-1} (x^2 + b^2)^\nu dx = b^{2\nu} a^{\lambda+\mu-1} B(\lambda, \mu) \times {}_3F_2 \left(-\nu, \frac{\lambda}{2}, \frac{\lambda+1}{2}; \frac{\lambda+\mu}{2}, \frac{\lambda+\mu+1}{2}; -\frac{a^2}{b^2} \right).$$

The conditions given in [3] are $\operatorname{Re}(\frac{a}{b}) > 0, \lambda > 0, \operatorname{Re} \mu > 0$. This entry appears as entry 186(10) of [2] as an example of the Riemann-Liouville transform

$$(5.2) \quad f(x) \mapsto \frac{1}{\Gamma(\mu)} \int_0^y f(x)(y-x)^{\mu-1} dx.$$

It is convenient to scale the formula, by the change of variables $x = at$, to the form

$$\int_0^1 t^{\lambda-1} (1-t)^{\mu-1} (1+c^2t^2)^\nu dt = B(\lambda, \mu) {}_3F_2 \left(-\nu, \frac{\lambda}{2}, \frac{\lambda+1}{2}; \frac{\lambda+\mu}{2}, \frac{\lambda+\mu+1}{2}; -c^2 \right),$$

with $c = a/b$. The binomial theorem gives

$$(5.3) \quad (1+c^2t^2)^\nu = {}_1F_0(-\nu; -; -c^2t^2) = \sum_{n=0}^{\infty} \frac{(-\nu)_n}{n!} (-1)^n c^{2n} t^{2n}$$

that produces

$$\begin{aligned} \int_0^1 t^{\lambda-1} (1-t)^{\mu-1} (1+c^2t^2)^\nu dt &= \sum_{n=0}^{\infty} \frac{(-\nu)_n}{n!} (-c^2)^n \int_0^1 t^{\lambda+2n-1} (1-t)^{\mu-1} dt \\ &= \sum_{n=0}^{\infty} \frac{(-\nu)_n}{n!} (-c^2)^n B(\lambda+2n, \mu). \end{aligned}$$

Now write the beta term as

$$\begin{aligned} B(\lambda+2n, \mu) &= \frac{\Gamma(\lambda+2n)\Gamma(\mu)}{\Gamma(\lambda+2n+\mu)} \\ &= \Gamma(\mu) \frac{2^{\lambda+2n-1} \Gamma(\frac{\lambda}{2}+n) \Gamma(\frac{\lambda+1}{2}+n)}{2^{\lambda+2n+\mu-1} \Gamma(\frac{\lambda+\mu}{2}+n) \Gamma(\frac{\lambda+\mu+1}{2}+n)} \end{aligned}$$

where the duplication formula for the gamma function

$$(5.4) \quad \Gamma(2x) = \frac{2^{2x-1}}{\sqrt{\pi}} \Gamma(x) \Gamma(x + \frac{1}{2})$$

has been employed. The relation $\Gamma(x+m) = (x)_m \Gamma(x)$ now yields

$$\int_0^1 t^{\lambda-1} (1-t)^{\mu-1} (1+c^2 t^2)^\nu dt = \frac{\Gamma(\mu)\Gamma(\frac{\lambda}{2})\Gamma(\frac{\lambda+1}{2})}{2^\mu \Gamma(\frac{\lambda+\mu}{2})\Gamma(\frac{\lambda+\mu+1}{2})} {}_3F_2 \left(-\nu, \frac{\lambda}{2}, \frac{\lambda+1}{2}; \frac{\lambda+\mu}{2}, \frac{\lambda+\mu+1}{2}; -c^2 \right).$$

Now simplify the gamma factors to produce the result.

Example 5.2. The next entry contains a typo in the 7th-edition of [3]. The correct version of **3.254.2** states that

$$(5.5) \quad \int_a^\infty x^{-\lambda} (x-a)^{\mu-1} (x^2+b^2)^\nu dx = a^{\mu-\lambda+2\nu} B(\mu, \lambda-\mu-2\nu) {}_3F_2 \left(-\nu, \frac{\lambda-\mu}{2} - \nu, \frac{1+\lambda-\mu}{2} - \nu; \frac{\lambda}{2} - \nu, \frac{1+\lambda}{2} - \nu; -\frac{b^2}{a^2} \right)$$

that follows directly from Example 5.1 by the change of variables $y = a^2/x$. It is convenient to scale this entry to the form

$$(5.6) \quad \int_1^\infty t^{-\lambda} (t-1)^{\mu-1} (t^2+c^2)^\nu dt = B(\mu, \lambda-\mu-2\nu) {}_3F_2 \left(-\nu, \frac{\lambda-\nu}{2} - \nu, \frac{1+\lambda-\mu}{2} - \nu; \frac{\lambda}{2} - \nu, \frac{1+\lambda}{2} - \nu; -c^2 \right).$$

6. A single factor of higher degree

In this section we consider entries in [3] of the

$$(6.1) \quad I = \int_a^b H_1(x)^{\mu-1} L_2(x)^{\nu-1} L_3(x)^{\lambda-1} dx$$

where $H_1(x)$ is a polynomial of degree $h \geq 2$ and L_j are linear functions.

Example 6.1. Entry **3.259.2** of [3] states that

$$\int_0^a x^{\nu-1} (a-x)^{\mu-1} (x^m+b^m)^\lambda dx = b^{m\lambda} a^{\mu+\nu-1} B(\mu, \nu) \times_{m+1} F_m \left(-\lambda, \frac{\nu}{m}, \frac{\nu+1}{m}, \dots, \frac{\nu+m-1}{m}; \frac{\mu+\nu}{m}, \frac{\mu+\nu+1}{m}, \dots, \frac{\mu+\nu+m-1}{m}; -\frac{a^m}{b^m} \right).$$

The scaling $t = x/a$ transforms this entry into

$$\int_0^1 t^{\nu-1} (1-t)^{\mu-1} (1+c^m t^m)^\lambda dt = B(\mu, \nu) \times_{m+1} F_m \left(-\lambda, \frac{\nu}{m}, \frac{\nu+1}{m}, \dots, \frac{\nu+m-1}{m}; \frac{\mu+\nu}{m}, \frac{\mu+\nu+1}{m}, \dots, \frac{\mu+\nu+m-1}{m}; -c^m \right)$$

with $c = a/b$. This is established next using a technique developed by Euler in his proof of the integral representation of ${}_2F_1$.

Start with

$$\begin{aligned} I &= \int_0^1 t^{\nu-1}(1-t)^{\mu-1}(c^m t^m + 1)^\lambda dt \\ &= \int_0^1 t^{\nu-1}(1-t)^{\mu-1} {}_1F_0(-\lambda; -; -c^m t^m) dt \end{aligned}$$

using the elementary identity (1.3). This gives

$$\begin{aligned} I &= \int_0^1 t^{\nu-1}(1-t)^{\mu-1} \sum_{n=0}^{\infty} \frac{(-\lambda)_n}{n!} (-c^m t^m)^n dt \\ &= \sum_{n=0}^{\infty} \frac{(-\lambda)_n}{n!} (-c^m)^n \int_0^1 t^{\nu+mn-1}(1-t)^{\mu-1} dt. \end{aligned}$$

The integral is recognized as a beta function value, therefore

$$\begin{aligned} I &= \sum_{n=0}^{\infty} \frac{(-\lambda)_n}{n!} (-c^m)^n \frac{\Gamma(\nu + mn)\Gamma(\mu)}{\Gamma(\nu + mn + \mu)} \\ &= \sum_{n=0}^{\infty} \frac{(-\lambda)_n}{n!} (-c^m)^n \frac{\Gamma(m(\frac{\nu}{m} + n))\Gamma(\mu)}{\Gamma(m(\frac{\nu+\mu}{m} + n))} \\ &= \Gamma(\mu) \sum_{n=0}^{\infty} \frac{(-\lambda)_n (-c^m)^n}{n!} \frac{m^{m(\nu/m+n)-1/2} \Gamma(\frac{\nu}{m} + n) \cdots \Gamma(\frac{\nu+m-1}{m} + n)}{m^{m(\frac{\nu+\mu}{m} + n)-1/2} \Gamma(\frac{\nu+\mu}{m} + n) \cdots \Gamma(\frac{\nu+\mu+m-1}{m} + n)} \\ &= \frac{\Gamma(\mu)}{m^\mu} \frac{\Gamma(\frac{\nu}{m}) \cdots \Gamma(\frac{\nu+m-1}{m})}{\Gamma(\frac{\nu+\mu}{m}) \cdots \Gamma(\frac{\nu+\mu+m-1}{m})} \times \sum_{n=0}^{\infty} \frac{(-\lambda)_n (\frac{\nu}{m})_n \cdots (\frac{\nu+m-1}{m})_n (-c^m)^n}{(\frac{\nu+\mu}{m})_n \cdots (\frac{\nu+\mu+m-1}{m})_n n!} \\ &= \frac{\Gamma(\mu)}{m^\mu} \frac{\Gamma(\frac{\nu}{m}) \cdots \Gamma(\frac{\nu+m-1}{m})}{\Gamma(\frac{\nu+\mu}{m}) \cdots \Gamma(\frac{\nu+\mu+m-1}{m})} \times \\ &\quad \times {}_{m+1}F_m\left(-\lambda, \frac{\nu}{m}, \dots, \frac{\nu+m-1}{m}; \frac{\nu+\mu}{m}, \dots, \frac{\nu+\mu+m-1}{m}; -c^m\right). \end{aligned}$$

This is the evaluation presented in entry **3.259.2**.

7. Integrals over a half-line

This section considers integrals over a half-line that can be expressed in terms of the hypergeometric function.

Example 7.1. To write (3.3) as an integral over an infinite half-line, make the change of variables $w = 1/y$ to obtain

$$(7.1) \quad \int_{1/u}^{\infty} w^{a-b-1}(1+w/r)^{-a} dw = \frac{u^b r^a}{b} {}_2F_1\left(a, b; b+1; -ru\right),$$

Now replace u by $1/u$ and r by $1/r$ to produce

$$(7.2) \quad \int_u^{\infty} w^{a-b-1}(1+rw)^{-a} dw = \frac{1}{bu^b r^a} {}_2F_1\left(a, b; b+1; -\frac{1}{ru}\right).$$

Finally let $b = a - s$ to obtain

$$(7.3) \quad \int_u^\infty w^{s-1}(1+rw)^{-a} dw = \frac{1}{(a-s)u^{a-s}r^a} {}_2F_1\left(a, a-s; a-s+1; -\frac{1}{ru}\right).$$

This appears as **3.194.2** in [3].

Example 7.2. The change of variable $y = 1/t$ converts (2.3) into **3.197.6**:

$$(7.4) \quad \int_1^\infty y^{a-c}(y-1)^{c-b-1}(\alpha y-1)^{-a} dy = \alpha^{-a}B(b, c-b) {}_2F_1(a, b; c; 1/\alpha)$$

where we have labelled $\alpha = 1/x$.

Example 7.3. The change of variables $y = t/(1-t)$ converts (2.3) into **3.197.5**:

$$(7.5) \quad \int_0^\infty y^{b-1}(1+y)^{a-c}(1+\alpha y)^{-a} dy = B(b, c-b) {}_2F_1(a, b; c; 1-\alpha)$$

where we have labelled $\alpha = 1-x$. If we now replace α by $1/\alpha$ we obtain

$$(7.6) \quad \int_0^\infty y^{b-1}(1+y)^{a-c}(y+\alpha)^{-a} dy = \alpha^a B(b, c-b) {}_2F_1(a, b; c; 1-1/\alpha).$$

Use the identity

$$(7.7) \quad {}_2F_1(a, b; c; 1-1/\alpha) = (1-\alpha)^a {}_2F_1(a, c-b; c; \alpha)$$

to produce **3.197.9**:

$$(7.8) \quad \int_0^\infty y^{b-1}(1+y)^{a-c}(y+\alpha)^{-a} dy = \alpha^a B(b, c-b) {}_2F_1(a, c-b; c; 1-\alpha).$$

Example 7.4. The change of variables $y = tu$ converts (2.3), with $-x$ instead of x , into **3.197.8**:

$$(7.9) \quad \int_0^u y^{b-1}(u-y)^{c-b-1}(y+\alpha)^{-a} dy = \alpha^{-a}u^{c-1}B(b, c-b) {}_2F_1(a, b; c; -u/\alpha)$$

where we have labelled $\alpha = u/x$.

Example 7.5. The change of variables $y = st/(1-t)$ converts (2.3) into

$$(7.10) \quad \int_0^\infty y^{b-1}(y+s)^{a-c}(y+r)^{-a} dy = r^{-a}s^{a+b-c}B(b, c-b) {}_2F_1\left(a, b; c; 1-\frac{s}{r}\right),$$

where $r = s/(1-x)$. This is **3.197.1** in [3]. The special case $a = c-1$ produces **3.227.1**:

$$(7.11) \quad \int_0^\infty \frac{y^{b-1}(y+r)^{1-c}}{y+s} dy = r^{1-c}s^{b-1}B(b, c-b) {}_2F_1\left(c-1, b; c; 1-\frac{s}{r}\right).$$

Example 7.6. Now shift the lower limit of integration via $x = y + u$ to produce

$$\int_u^\infty (x-u)^{b-1}(x-u+s)^{a-c}(x-u+r)^{-a} dx = r^{-a}u^{a+b-c}B(b, c-b) {}_2F_1\left(a, b; c; 1-\frac{s}{r}\right).$$

Choose $s = u$ and introduce the parameter v by $v = r - u$ to get

$$\int_u^\infty x^{a-c}(x-u)^{b-1}(x+v)^{-a} dx = (v+u)^{-a} u^{a+b-c} B(b, c-b) {}_2F_1\left(a, b; c; \frac{v}{v+u}\right).$$

Introduce new parameters via $a = -p$, keeping b and $c = q - p$. This yields

$$\begin{aligned} \int_u^\infty x^{-q}(x-u)^{b-1}(x+v)^p dx &= (v+u)^p u^{b-q} B(b, c-b-p) {}_2F_1\left(-p, b; q-p; \frac{v}{v+u}\right) \\ &= (v+u)^p u^{b-q} B(b, c-b-p) {}_2F_1\left(b, -p; q-p; \frac{v}{v+u}\right) \end{aligned}$$

where the symmetry of the hypergeometric function in its two variables has been used.

This result is transformed using **9.131.1**:

$$(7.12) \quad {}_2F_1(a, b; c; z) = (1-z)^{-a} {}_2F_1(a, c-b; c; z/(z-1)),$$

that gives

$$\int_u^\infty x^{-q}(x-u)^{b-1}(x+v)^p dx = (v+u)^{b+p} u^{b-q} B(b, q-p-b) {}_2F_1\left(b, q; q-p; -\frac{v}{u}\right).$$

This is the form that is found in **3.197.2**.

8. An exponential scale

The change of variables $t = e^{-r}$ in (2.3) produces

$$(8.1) \quad {}_2F_1(a, b; c; x) = \frac{1}{B(b, c-b)} \int_0^\infty e^{-br} (1-e^{-r})^{c-b-1} (1-xe^{-r})^{-a} dr.$$

The parameters are relabeled by $a = \rho$, $b = \mu$, $c = \nu + \mu$, $x = \beta$ to produce **3.312.3**:

$$(8.2) \quad \int_0^\infty (1-e^{-x})^{\nu-1} (1-\beta e^{-x})^{-\rho} e^{-\mu x} dx = B(\mu, \nu) {}_2F_1(\rho, \mu; \mu + \nu; \beta).$$

9. A more challenging example

The evaluation of **3.197.7**

$$(9.1) \quad \int_0^\infty x^{\mu-1/2}(x+s)^{-\mu}(x+r)^{-\mu} dx = \sqrt{\pi}(\sqrt{r} + \sqrt{s})^{1-2\mu} \frac{\Gamma(\mu-1/2)}{\Gamma(\mu)}$$

requires some more properties of the hypergeometric function.

The scaling $x = rt$ produces

$$(9.2) \quad I = s^{-\mu} \sqrt{r} \int_0^\infty t^{\mu-1/2} (1+t)^{-\mu} (1+rt/s)^\mu dt$$

and using **3.197.5** we have

$$(9.3) \quad I = s^{-\mu} \sqrt{r} B\left(\mu + \frac{1}{2}, \mu - \frac{1}{2}\right) {}_2F_1\left(\mu, \mu + \frac{1}{2}, 2\mu; z\right)$$

where $z = 1 - r/s$. To simplify this expression we employ the relation

$$\begin{aligned} {}_2F_1(\alpha, \beta; \gamma; z) &= \frac{(1-z)^{-\alpha} \Gamma(\gamma) \Gamma(\beta - \alpha)}{\Gamma(\beta) \Gamma(\gamma - \alpha)} {}_2F_1(\alpha, \gamma - \beta; \alpha - \beta + 1; \frac{1}{1-z}) + \\ &+ \frac{(1-z)^{-\beta} \Gamma(\gamma) \Gamma(\alpha - \beta)}{\Gamma(\beta) \Gamma(\gamma - \beta)} {}_2F_1(\beta, \gamma - \alpha; \beta - \alpha + 1; \frac{1}{1-z}) \end{aligned}$$

to produce

$$\begin{aligned} {}_2F_1\left(\mu, \mu + \frac{1}{2}, 2\mu; z\right) &= \frac{(1-z)^{-\mu} \Gamma(2\mu) \Gamma(1/2)}{\Gamma(\mu + 1/2) \Gamma(\mu)} {}_2F_1\left(\mu, \mu - \frac{1}{2}; \frac{1}{2}; \frac{1}{1-z}\right) \\ &+ \frac{(1-z)^{-\mu-1/2} \Gamma(2\mu) \Gamma(-1/2)}{\Gamma(\mu - 1/2) \Gamma(\mu)} {}_2F_1\left(\mu, \mu + \frac{1}{2}; \frac{3}{2}; \frac{1}{1-z}\right). \end{aligned}$$

The binomial theorem shows that

$$(9.4) \quad {}_2F_1\left(-\frac{n}{2}, -\frac{n-1}{2}; \frac{1}{2}; \frac{z^2}{t^2}\right) = \frac{1}{2t^n} ((t+z)^n + (t-z)^n),$$

that appears as **9.121.2** in [3]. Thus

$${}_2F_1\left(\mu, \mu - \frac{1}{2}; \frac{1}{2}; \frac{1}{1-z}\right) = \frac{1}{2(1-z)^{1/2-\mu}} ((1 + \sqrt{1-z})^{1-2\mu} + (-1 + \sqrt{1-z})^{1-2\mu}).$$

Similarly, **9.121.4** states that

$$(9.5) \quad {}_2F_1\left(-\frac{n-1}{2}, -\frac{n-2}{2}; \frac{3}{2}; \frac{z^2}{t^2}\right) = \frac{1}{2nzt^{n-1}} ((t+z)^n - (t-z)^n),$$

to produce

$${}_2F_1\left(\mu, \mu - \frac{1}{2}; \frac{3}{2}; \frac{1}{1-z}\right) = \frac{1}{2(1-2\mu)(1-z)^{-\mu}} ((1 + \sqrt{1-z})^{1-2\mu} - (-1 + \sqrt{1-z})^{1-2\mu}).$$

Replacing these values in (9.3) produces the result.

10. One last example: a combination of algebraic factors and exponentials

Entry **3.389.1** presents an analytic expression for the integral

$$(10.1) \quad I := \int_0^a x^{2\nu-1} (a^2 - x^2)^{\rho-1} e^{\mu x} dx.$$

The evaluation begins with an elementary scaling to obtain

$$\begin{aligned} I &= a^{2(\rho-1)} \int_0^1 x^{2\nu-1} \left(1 - \frac{x^2}{a^2}\right)^{\rho-1} e^{\mu x} dx \\ &= \frac{1}{2} a^{2\rho-1} \int_0^1 (ay^{1/2})^{2\nu-1} (1-y)^{\rho-1} e^{\mu ay^{1/2}} y^{-1/2} dy. \end{aligned}$$

Now use ${}_0F_0(; ; x) = e^x$ to obtain

$$\begin{aligned} I &= \frac{a^{2\rho+2\nu-2}}{2} \int_0^1 y^{\nu-1} (1-y)^{\rho-1} {}_0F_0(; ; \mu a y^{1/2}) dy \\ &= \frac{a^{2\rho+2\nu-2}}{2} \int_0^1 y^{\nu-1} (1-y)^{\rho-1} \sum_{n=0}^{\infty} \frac{(\mu a y^{1/2})^n}{n!} dy \\ &= \frac{a^{2\rho+2\nu-2}}{2} \sum_{n=0}^{\infty} \frac{(\mu a)^n}{n!} \int_0^1 y^{\nu+n/2-1} (1-y)^{\rho-1} dy. \end{aligned}$$

The integral is now recognized as a beta value to conclude that

$$\begin{aligned} I &= \frac{a^{2\rho+2\nu-2}}{2} \sum_{n=0}^{\infty} \frac{(\mu a)^n}{n!} B(\nu + n/2, \rho) \\ &= \frac{a^{2\rho+2\nu-2} \Gamma(\rho)}{2} \sum_{n=0}^{\infty} \frac{(\mu a)^n}{n!} \frac{\Gamma(\nu + n/2)}{\Gamma(\nu + n/2 + \rho)} \\ &= \frac{a^{2\rho+2\nu-2} \Gamma(\rho) \Gamma(\nu)}{2\Gamma(\nu + \rho)} \sum_{k=0}^{\infty} \frac{(\mu a)^{2k} (\nu)_k}{\Gamma(2k+1) (\nu + \rho)_k} + \frac{a^{2\rho+2\nu-2} \Gamma(\rho)}{2} \sum_{k=0}^{\infty} \frac{(\mu a)^{2k+1} \Gamma(\nu + k + 1/2)}{(2k+1)! \Gamma(\nu + \rho + k + 1/2)} \end{aligned}$$

and combining the gamma factors to produce the beta function yields

$$\begin{aligned} I &= \frac{1}{2} a^{2\rho+2\nu-2} B(\rho, \nu) \sum_{k=0}^{\infty} \frac{(\mu^2 a^2)^k (\nu)_k}{(2k) \Gamma(2k) (\nu + \rho)_k} + \\ &\quad + \frac{1}{2} a^{2\rho+2\nu-1} \mu \Gamma(\rho) \sum_{k=0}^{\infty} \frac{(\mu a)^{2k}}{\Gamma(2k+2)} \frac{(\nu + 1/2)_k \Gamma(\nu + 1/2)}{(\nu + \rho + 1/2)_k \Gamma(\nu + \rho + 1/2)}. \end{aligned}$$

This can be reduced to

$$\begin{aligned} 2I &= a^{2\rho+2\nu-2} B(\rho, \nu) \sum_{k=0}^{\infty} \frac{(\nu)_k (\mu^2 a^2)^k}{(\nu + \rho)_k (2k)} \frac{2^{1-2k} \sqrt{\pi}}{\Gamma(k) \Gamma(k + 1/2)} + \\ &\quad + a^{2\rho+2\nu-1} \mu B(\rho, \nu + 1/2) \sum_{k=0}^{\infty} \frac{(\nu + 1/2)_k}{(\nu + \rho + 1/2)_k} \frac{(\mu^2 a^2)^k 2^{1-2(k+1)} \sqrt{\pi}}{\Gamma(k+1) \Gamma(k + \frac{3}{2})} \\ &= a^{2\rho+2\nu-2} B(\rho, \nu) \sum_{k=0}^{\infty} \frac{(\nu)_k}{(\nu + \rho)_k (\frac{1}{2})_k k!} \left(\frac{\mu^2 a^2}{4} \right)^k + \\ &\quad + a^{2\rho+2\nu-1} \mu B(\rho, \nu + 1/2) \sum_{k=0}^{\infty} \frac{(\nu + 1/2)_k}{(\nu + \rho + 1/2)_k (\frac{3}{2})_k} \left(\frac{\mu^2 a^2}{4} \right)^k \\ &= a^{2\rho+2\nu-2} B(\rho, \nu) {}_1F_2 \left(\nu; \nu + \rho, \frac{1}{2}; \frac{\mu^2 a^2}{4} \right) + \\ &\quad + a^{2\rho+2\nu-1} \mu B(\rho, \nu + 1/2) {}_1F_2 \left(\nu + 1; \nu + \rho + 1/2, \frac{3}{2}; \frac{\mu^2 a^2}{4} \right). \end{aligned}$$

There are many other entries of [3] that can be evaluated in terms of hypergeometric functions. A second selection of examples is in preparation.

References

- [1] G. Andrews, R. Askey, and R. Roy. *Special Functions*, volume 71 of *Encyclopedia of Mathematics and its Applications*. Cambridge University Press, New York, 1999.
- [2] A. Erdélyi. *Tables of Integral Transforms*, volume II. McGraw-Hill, New York, 1st edition, 1954.
- [3] I. S. Gradshteyn and I. M. Ryzhik. *Table of Integrals, Series, and Products*. Edited by A. Jeffrey and D. Zwillinger. Academic Press, New York, 7th edition, 2007.

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