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# The integrals in Gradshteyn and Ryzhik. <br> Part 27: More logarithmic examples 

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#### Abstract

The table of Gradshteyn and Ryzhik contains many entries where the integrand is a combination of an elementary function and the logarithmic of another function of the same type. This paper presents proofs of some of these. A sample of examples where the elementary function is replaced by an algebraic function is also discussed.


## 1. Introduction

The compendium [5] contains a large collection of evaluation of integrals of the form

$$
\begin{equation*}
\int_{a}^{b} R_{1}(x) \ln R_{2}(x) d x \tag{1.1}
\end{equation*}
$$

where $R_{1}$ and $R_{2}$ are rational functions. The first paper in this series [ $\left.\mathbf{9}\right]$ considered the family

$$
\begin{equation*}
f_{n}(a)=\int_{0}^{\infty} \frac{\ln ^{n-1} x d x}{(x-1)(x+a)}, \text { for } n \geqslant 2 \text { and } a>0 \tag{1.2}
\end{equation*}
$$

The function $f_{n}(a)$ is given explicitly by

$$
\begin{align*}
f_{n}(a) & =\frac{(-1)^{n}(n-1)!}{1+a}\left[1+(-1)^{n}\right] \zeta(n)  \tag{1.3}\\
& +\frac{1}{n(1+a)} \sum_{j=0}^{\lfloor n / 2\rfloor}\binom{n}{2 j}\left(2^{2 j}-2\right)(-1)^{j-1} B_{2 j} \pi^{2 j}(\log a)^{n-2 j} .
\end{align*}
$$

Here $\zeta(s)$ is the Riemann zeta function and $B_{2 j}$ is the Bernoulli number. In particular, (1.3) shows that $(1+a) f_{n}(a)$ is a polynomial in $\log a$.

[^0]Other papers in this series $[\mathbf{3}, \mathbf{8}, \mathbf{1 0}]$ and also $[\mathbf{6}]$ considered examples of integrals of this type. The results in [3] can be used to provide explicit expressions for an integral of the type considered here, when the poles of the rational function $R_{2}$ in (1.1) have real or purely imaginary parts. The present paper is a continuation of this work.

## 2. Some examples involving rational functions

This section considers of integrals of the form

$$
\begin{equation*}
\int_{a}^{b} R_{1}(x) \ln R_{2}(x) d x \tag{2.1}
\end{equation*}
$$

where $R_{1}$ and $R_{2}$ are rational functions.
Example 2.1. Entry 4.234.4 is

$$
\begin{equation*}
\int_{0}^{\infty} \frac{1-x^{2}}{\left(1+x^{2}\right)^{2}} \ln x d x=-\frac{\pi}{2} \tag{2.2}
\end{equation*}
$$

To evaluate this entry, observe that

$$
\begin{equation*}
\frac{d}{d x} \frac{x}{1+x^{2}}=\frac{1-x^{2}}{\left(1+x^{2}\right)^{2}} \tag{2.3}
\end{equation*}
$$

and integrating by parts gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{1-x^{2}}{\left(1+x^{2}\right)^{2}} \ln x d x=-\int_{0}^{\infty} \frac{d x}{1+x^{2}}=-\frac{\pi}{2} \tag{2.4}
\end{equation*}
$$

Example 2.2. Entry 4.234 .5 states that

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{2} \ln x d x}{\left(1-x^{2}\right)\left(1+x^{4}\right)}=-\frac{\pi^{2}}{16(2+\sqrt{2})} \tag{2.5}
\end{equation*}
$$

To prove this use the method of partial fraction to obtain

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{2} \ln x d x}{\left(1-x^{2}\right)\left(1+x^{4}\right)}=\frac{1}{4} \int_{0}^{1} \frac{\ln x d x}{1-x}+\frac{1}{4} \int_{0}^{1} \frac{\ln x d x}{1+x}+\frac{1}{2} \int_{0}^{1} \frac{\left(x^{2}-1\right) \ln x d x}{1+x^{4}} \tag{2.6}
\end{equation*}
$$

The first integral is $-\pi^{2} / 6$ according to entry 4.231 .2 and the second one is $-\pi^{2} / 12$ from entry 4.231 .1 . These entries were established in $[\mathbf{1}]$. This gives

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{2} \ln x d x}{\left(1-x^{2}\right)\left(1+x^{4}\right)}=-\frac{\pi^{2}}{16}+\frac{1}{2} \int_{0}^{1} \frac{\left(x^{2}-1\right) \ln x d x}{1+x^{4}} \tag{2.7}
\end{equation*}
$$

To evaluate the last integral, observe that

$$
\begin{equation*}
\frac{x^{2}-1}{1+x^{4}}=\sum_{n=0}^{\infty}(-1)^{n-1} x^{4 n}+\sum_{n=0}^{\infty}(-1)^{n} x^{4 n+2} \tag{2.8}
\end{equation*}
$$

Now recall the digamma function $\psi(z)=\Gamma^{\prime}(z) / \Gamma(z)$ and the expansion of its derivative

$$
\begin{equation*}
\psi^{\prime}(x)=\sum_{n=0}^{\infty} \frac{1}{(x+n)^{2}} \tag{2.9}
\end{equation*}
$$

Details about this function may be found in [4] and [13]. This gives

$$
\begin{equation*}
\int_{0}^{1} \frac{\left(x^{2}-1\right) \ln x d x}{1+x^{4}}=\frac{1}{64}\left[\psi^{\prime}\left(\frac{1}{8}\right)-\psi^{\prime}\left(\frac{3}{8}\right)-\psi^{\prime}\left(\frac{5}{8}\right)+\psi^{\prime}\left(\frac{7}{8}\right)\right] . \tag{2.10}
\end{equation*}
$$

The classical relation

$$
\begin{equation*}
\Gamma(x) \Gamma(1-x)=\frac{\pi}{\sin \pi x} \tag{2.11}
\end{equation*}
$$

can be shifted to produce

$$
\begin{equation*}
\Gamma\left(\frac{1}{2}+x\right) \Gamma\left(\frac{1}{2}-x\right)=\frac{\pi}{\cos \pi x} \tag{2.12}
\end{equation*}
$$

Logarithmic differentiation shows that the digamma function satisfies

$$
\begin{equation*}
\psi\left(\frac{1}{2}+x\right)-\psi\left(\frac{1}{2}-x\right)=\pi \tan \pi x \tag{2.13}
\end{equation*}
$$

This appears as Entry 8.365.9 in [5]. Differentiation produces

$$
\begin{equation*}
\psi^{\prime}\left(\frac{1}{2}+x\right)+\psi^{\prime}\left(\frac{1}{2}-x\right)=\pi^{2} \sec ^{2} \pi x . \tag{2.14}
\end{equation*}
$$

Now use (2.14) and group $1 / 8$ with $7 / 8$ and $3 / 8$ with $5 / 8$ to produce

$$
\begin{equation*}
\int_{0}^{1} \frac{\left(x^{2}-1\right) \ln x d x}{1+x^{4}}=\frac{1}{64}\left(\frac{4 \pi^{2}}{2-\sqrt{2}}-\frac{4 \pi^{2}}{2+\sqrt{2}}\right)=\frac{\pi^{2}}{8 \sqrt{2}} \tag{2.15}
\end{equation*}
$$

Note 2.3. The reader should evaluate the family of integrals

$$
\begin{equation*}
I_{n}=\int_{0}^{1} \frac{x^{2 n} \ln x}{\left(1-x^{2}\right)\left(1+x^{4}\right)^{n}} d x, \quad n \in \mathbb{N}, \tag{2.16}
\end{equation*}
$$

by the method described here. The computation of the first few special values indicates an interesting arithmetic structure of the answer.

## 3. An entry involving the Poisson kernel for the disk

The section discusses a single entry in [5], where the integrand involves the Poisson kernel for the disk. Further examples of this type will be presented in a future publication.

Example 3.1. The next evaluation is Entry 4.233.5:

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\ln x d x}{x^{2}+2 x a \cos t+a^{2}}=\frac{t}{\sin t} \frac{\ln a}{a} \tag{3.1}
\end{equation*}
$$

The integrand is related to the Poisson kernel for the unit disk $D=\{z \in \mathbb{C}:|z|<1\}$.
Theorem 3.2. Define

$$
\begin{equation*}
\mathcal{P}_{r}(\theta)=\operatorname{Re} \frac{1+r e^{i \theta}}{1-r e^{i \theta}} \tag{3.2}
\end{equation*}
$$

then $\mathcal{P}_{r}(\theta)$ is given by

$$
\begin{equation*}
\mathcal{P}_{r}(\theta)=\sum_{n=-\infty}^{\infty} r^{|n|} e^{i n \theta}=\frac{1-r^{2}}{1-2 r \cos \theta+r^{2}} \tag{3.3}
\end{equation*}
$$

Moreover, given $f$ defined on the boundary of $D$, the expression

$$
\begin{equation*}
u\left(r e^{i \theta}\right)=\frac{1}{2 \pi} \int_{-\pi}^{\pi} \mathcal{P}_{r}(\theta-t) f\left(e^{i t}\right) d t \tag{3.4}
\end{equation*}
$$

for $0 \leqslant r<1$, is a harmonic function on $D$ and it has a radial limit which agrees with $f$ almost everywhere on the boundary of $D$.

The form of the Poisson kernel can be used to establish the next result.
Lemma 3.3. For $a, x \in \mathbb{R}$ with $|x|<|a|$,

$$
\begin{equation*}
\sum_{k=0}^{\infty} \frac{(-1)^{k} \sin ((k+1) t) x^{k}}{a^{k}}=\frac{a^{2} \sin t}{x^{2}+2 a x \cos t+a^{2}} \tag{3.5}
\end{equation*}
$$

Note 3.4. The Chebyshev polynomial of the second kind $U_{n}(t)$ is defined by the identity

$$
\begin{equation*}
\frac{\sin ((n+1) \theta)}{\sin \theta}=U_{n}(\cos \theta) \tag{3.6}
\end{equation*}
$$

The result of Lemma 3.3 can be written as

$$
\begin{equation*}
\sum_{k=0}^{\infty} U_{k}(t) x^{k}=\frac{1}{x^{2}-2 x \cos t+1} \tag{3.7}
\end{equation*}
$$

Lemma 3.3 produces

$$
\begin{equation*}
\int_{0}^{R} \frac{x^{s} d x}{x^{2}+2 a x \cos t+a^{2}}=\frac{1}{a^{2} \sin t} \sum_{k=0}^{\infty} \frac{(-1)^{k} \sin ((k+1) t) R^{k+s+1}}{a^{k}(k+s+1)} \tag{3.8}
\end{equation*}
$$

Now write $\sin ((k+1) t)$ in terms of exponential to obtain an expression for the previous integral as
$\int_{0}^{R} \frac{x^{s} d x}{x^{2}+2 a x \cos t+a^{2}}=\frac{R^{s+1}}{2 i a^{2} \sin t}\left(e^{i t} \Phi\left(-\frac{R}{a e^{i t}}, 1, s+1\right)-e^{-i t} \Phi\left(-\frac{R}{a e^{-i t}}, 1, s+1\right)\right)$
where

$$
\begin{equation*}
\Phi(z, s, a)=\sum_{k=0}^{\infty} \frac{z^{k}}{(a+k)^{s}} \tag{3.9}
\end{equation*}
$$

is the Lerch Phi function.
Now differentiate with respect to $s$ and let $s \rightarrow 0$ to produce
(3.10) $\int_{0}^{R} \frac{\ln x d x}{x^{2}+2 a x \cos t+a^{2}}=\frac{i \ln R}{2 a \sin t}\left(\log \left(1+e^{-i t} R / a\right)-\log \left(1+e^{i t} R / a\right)\right)$

$$
+\frac{i}{2 a \sin t}\left(\operatorname{Li}_{2}\left(-e^{-i t} R / a\right)-\operatorname{Li}_{2}\left(-e^{-i t} R / a\right)\right),
$$

where

$$
\begin{equation*}
\mathrm{Li}_{2}(z)=\sum_{k=1}^{\infty} \frac{z^{k}}{k^{2}} \tag{3.11}
\end{equation*}
$$

is the dilogarithm function. Then use the identity

$$
\begin{equation*}
i\left(\operatorname{Li}_{2}\left(-e^{-i t} R / a\right)-\operatorname{Li}_{2}\left(-e^{-i t} R / a\right)\right)=-\int_{0}^{t} \ln \left(\frac{a^{2}+2 R a \cos z+R^{2}}{a^{2}}\right) d z \tag{3.12}
\end{equation*}
$$

to obtain

$$
\begin{aligned}
\text { (3.13) } \int_{0}^{R} \frac{\ln x d x}{x^{2}+2 a x \cos t+a^{2}}= & \frac{i \ln R}{2 a \sin t}\left(\log \left(1+e^{-i t} R / a\right)-\log \left(1+e^{i t} R / a\right)\right) \\
& -\frac{1}{2 a \sin t} \int_{0}^{t} \ln \left(\frac{a^{2}+2 R a \cos z+R^{2}}{a^{2}}\right) d z .
\end{aligned}
$$

The next step is to differentiate (3.13) with respect to $t$ and let $R \rightarrow \infty$. The left-hand side produces

$$
\begin{equation*}
T_{1}(a, t)=\int_{0}^{\infty} \frac{2 a x \ln x \sin t d x}{\left(x^{2}+2 a x \cos t+a^{2}\right)^{2}} \tag{3.14}
\end{equation*}
$$

Direct differentiation of the right-hand side yields

$$
\begin{equation*}
T_{2}(a, t)=\lim _{R \rightarrow \infty} V_{1}(R ; a, t)+V_{2}(R ; a, t) \tag{3.15}
\end{equation*}
$$

where
(3.16) $V_{1}(R ; a, t)=\frac{R \ln R(R+a \cos t)}{a \sin t\left(a^{2}+2 a R \cos t+R^{2}\right)}-\frac{1}{2 a \sin t} \ln \left(\frac{a^{2}+2 a R \cos t+R^{2}}{a^{2}}\right)$
and

$$
\begin{align*}
V_{2}(R ; a, t)= & \frac{i \cos t \ln R}{2 a \sin ^{2} t}\left(\log \left(1+e^{i t} R / a\right)-\log \left(1+e^{-i t} R / a\right)\right)  \tag{3.17}\\
& +\frac{\cos t}{2 a \sin ^{2} t} \int_{0}^{t} \ln \left(\frac{a^{2}+2 R a \cos z+R^{2}}{a^{2}}\right) d z .
\end{align*}
$$

Proposition 3.5. The function $T_{2}(a, t)$ is given

$$
\begin{equation*}
T_{2}(a, t)=-\frac{\ln a}{2 a \sin t}(t \cot t-1) . \tag{3.18}
\end{equation*}
$$

Proof. Start with the computation of the limiting behavior of $V_{1}(R ; a, t)$. The claim that

$$
\begin{equation*}
\lim _{R \rightarrow \infty} V_{1}(R ; a, t)=\frac{\ln a}{a \sin (t)} \tag{3.19}
\end{equation*}
$$

is verified first.
First note that since

$$
\begin{equation*}
\lim _{R \rightarrow \infty} \frac{R \ln R}{a^{2}+2 a R \cos (t)+R^{2}}=0 \tag{3.20}
\end{equation*}
$$

then
$\lim _{R \rightarrow \infty} V_{1}(R ; a, t)=\frac{1}{a \sin t} \lim _{R \rightarrow \infty}\left(\frac{R^{2} \ln R}{a^{2}+2 a R \cos t+R^{2}}-\frac{1}{2} \ln \left(a^{2}+2 a R \cos t+R^{2}\right)+\ln a\right)$.

The claim is equivalent to

$$
\begin{equation*}
\lim _{R \rightarrow \infty}\left(\frac{R^{2} \ln R}{a^{2}+2 a R \cos t+R^{2}}-\frac{1}{2} \ln \left(a^{2}+2 a R \cos t+R^{2}\right)\right)=0 . \tag{3.21}
\end{equation*}
$$

The identities

$$
\begin{equation*}
\frac{R^{2} \ln R}{a^{2}+2 a R \cos t+R^{2}}=\frac{\ln R}{a^{2} / R^{2}+2 a \cos t / R+1} \tag{3.22}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{1}{2} \ln \left(a^{2}+2 a R \cos t+R^{2}\right)=\ln R+\frac{1}{2} \ln \left(a^{2} / R^{2}+2 a \cos t / R+1\right) \tag{3.23}
\end{equation*}
$$

can be used to see that the left-hand side of (3.21) is equivalent to

$$
\lim _{R \rightarrow \infty}\left(\ln R\left(\frac{1}{a^{2} / R^{2}+2 a \cos t / R+1}-1\right)-\frac{1}{2} \ln \left(a^{2} / R^{2}+2 a \cos t / R+1\right)\right)=0
$$

It is clear that the second term vanishes as $R \rightarrow \infty$. For the first term, observe that

$$
\begin{equation*}
\frac{1}{a^{2} / R^{2}+2 a \cos (t) / R+1}-1=-\frac{2 a \cos t}{R}+O\left(\frac{1}{R^{2}}\right) \tag{3.24}
\end{equation*}
$$

and thus the first term also vanishes as $R \rightarrow \infty$. This concludes the proof.
The next step is to verify that

$$
\begin{align*}
V_{2}(R ; a, t)= & \frac{i \cot t \ln R}{2 a \sin ^{2} t}\left(\log \left(1+e^{i t} R / a\right)-\log \left(1+e^{-i t} R / a\right)\right)  \tag{3.25}\\
& +\frac{\cos t}{2 a \sin ^{2} t} \int_{0}^{t} \ln \left(\frac{a^{2}+2 a R \cos z+R^{2}}{a^{2}}\right) d z
\end{align*}
$$

satisfies

$$
\begin{equation*}
\lim _{R \rightarrow \infty} V_{2}(R ; a, t)=-\frac{t \cos t}{a \sin ^{2} t} \ln a \tag{3.26}
\end{equation*}
$$

The proof begins with the identity

$$
\begin{equation*}
\log (1+b / x)=\log (b / x)+\sum_{n=1}^{\infty}(-1)^{n-1} \frac{x^{n}}{n b^{n}} \tag{3.27}
\end{equation*}
$$

to obtain
(3.28) $\log \left(1+e^{i t} R / a\right)-\log \left(1+e^{-i t} R / a\right)=\log \left(e^{i t}\right)-\log \left(e^{-i t}\right)+O(a / R)$, as $R \rightarrow \infty$.

The bounds $0<t<\pi$ imply $\log \left(e^{i t}\right)-\log \left(e^{-i t}\right)=2 i t$. This gives

$$
\begin{aligned}
\lim _{R \rightarrow \infty} V_{2}(R ; a, t) & =\lim _{R \rightarrow \infty}\left(\frac{\cos t}{2 a \sin ^{2} t} \int_{0}^{t} \ln \left(\frac{a^{2}+2 a R \cos z+R^{2}}{a^{2}}\right) d z-\frac{t \cos z \ln R}{a \sin ^{2} t}\right) \\
& =\lim _{R \rightarrow \infty} \frac{\cos t}{2 a \sin ^{2} t}\left(\int_{0}^{t} \ln \left(\frac{a^{2}+2 a R \cos z+R^{2}}{a^{2}}\right) d z-2 t \ln R\right) \\
& =\lim _{R \rightarrow \infty} \frac{\cos t}{2 a \sin ^{2} t}\left(\int_{0}^{t} \ln \left(\frac{a^{2}+2 a R \cos z+R^{2}}{a^{2}}\right)-\ln \left(R^{2}\right) d z\right) \\
& =\lim _{R \rightarrow \infty} \frac{\cos t}{2 a \sin ^{2} t}\left(\int_{0}^{t}\left[\ln \left(a^{2}+2 a R \cos z+R^{2}\right)-\ln \left(R^{2}\right)\right] d z-2 t \ln a\right) .
\end{aligned}
$$

The identity

$$
\int_{0}^{t}\left[\ln \left(a^{2}+2 a R \cos z+R^{2}\right)-\ln \left(R^{2}\right)\right] d z=\int_{0}^{t} \ln \left(\frac{a^{2}}{R^{2}}+\frac{2 a \cos z}{R}+1\right) d z
$$

gives the result. The proof of the Proposition is finished.

The evaluation of entry 4.233 .5 is now obtained from the identity $T_{1}(a, t)=$ $T_{2}(a, t)$. Observe that this implies

$$
\begin{equation*}
\int_{0}^{\infty} \frac{2 a x \ln x \sin t d x}{\left(x^{2}+2 a x \cos t+a^{2}\right)^{2}}=-\frac{\ln a}{a \sin t}(t \cot t-1) . \tag{3.29}
\end{equation*}
$$

Integrating with respect to $t$ gives (3.1). Entry 4.231.8 in [5], established in [3],

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\ln x d x}{x^{2}+a^{2}}=\frac{\pi \ln a}{2 a} \tag{3.30}
\end{equation*}
$$

can be used to show that the implicit constant of integration actually vanishes. The evaluation is complete.

## 4. Some rational integrands with a pole at $x=1$

This section contains proofs of the four entries appearing in Section 4.235. These are integrals of the form

$$
\begin{equation*}
f(a, b, c):=\int_{0}^{\infty} \frac{x^{b}-x^{c}}{1-x^{a}} \ln x d x \tag{4.1}
\end{equation*}
$$

where $a, b, c \in \mathbb{N}$. These integrals are evaluated using entry 4.254 .2

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1} \ln x}{1-x^{q}} d x=-\frac{\pi^{2}}{q^{2} \sin ^{2} \frac{\pi p}{q}} \tag{4.2}
\end{equation*}
$$

To obtain this formula, start from 3.231.6

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1}-x^{q-1}}{1-x} d x=\pi(\cot \pi p-\cot \pi q) \tag{4.3}
\end{equation*}
$$

established in $[\mathbf{7}]$ and make the change of variables $t=x^{q}$ to produce

$$
\begin{aligned}
\int_{0}^{\infty} \frac{x^{p-1}-1}{1-x^{q}} d x & =-\frac{1}{q} \int_{0}^{\infty} \frac{t^{1 / q-1}-t^{p / q-1}}{1-t} d t \\
& =-\frac{\pi}{q}\left(\cot \frac{\pi}{q}-\cot \frac{\pi p}{q}\right)
\end{aligned}
$$

Differentiating with respect to $p$ gives (4.2).
Lemma 4.1. Let $a, b, c \in \mathbb{R}$. Then

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{b-1}-x^{c-1}}{1-x^{a}} \ln x d x=-\frac{\pi^{2}}{a^{2}} \frac{\sin \left(c_{1}-b_{1}\right) \sin \left(c_{1}+b_{1}\right)}{\sin ^{2} b_{1} \sin ^{2} c_{1}} \tag{4.4}
\end{equation*}
$$

where $b_{1}=\pi b / a$ and $c_{1}=\pi c / a$.
Proof. Simply write

$$
\int_{0}^{\infty} \frac{x^{b-1}-x^{c-1}}{1-x^{a}} \ln x d x=\int_{0}^{\infty} \frac{x^{b-1}}{1-x^{a}} \ln x d x-\int_{0}^{\infty} \frac{x^{c-1}}{1-x^{a}} \ln x d x
$$

and use (4.2).
The four entries in Section 4.235 are established next.
Example 4.1. Entry 4.235.1 states that

$$
\begin{equation*}
\int_{0}^{\infty} \frac{(1-x) x^{n-2}}{1-x^{2 n}} \ln x d x=-\frac{\pi^{2}}{4 n^{2}} \tan ^{2} \frac{\pi}{2 n} \tag{4.5}
\end{equation*}
$$

Lemma 4.1 is used with $a=2 n, b=n-1$ and $c=n$. This gives

$$
\begin{equation*}
b_{1}=\frac{\pi}{2}-\frac{\pi}{2 n} \text { and } c_{1}=\frac{\pi}{2} \tag{4.6}
\end{equation*}
$$

and

$$
\int_{0}^{\infty} \frac{(1-x) x^{n-2}}{1-x^{2 n}} \ln x d x=-\frac{\pi^{2}}{4 n^{2}} \frac{\sin \left(\frac{\pi}{2}-\frac{\pi}{2 n}\right) \sin \left(\frac{\pi}{2}+\frac{\pi}{2 n}\right)}{\sin ^{2}\left(\frac{\pi}{2}-\frac{\pi}{2 n}\right)}=-\frac{\pi^{2}}{4 n^{2}} \tan ^{2} \frac{\pi}{2 n}
$$

Example 4.2. Entry 4.235.2 is

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\left(1-x^{2}\right) x^{m-1}}{1-x^{2 n}} \ln x d x=-\frac{\pi^{2}}{4 n^{2}} \frac{\sin \left(\frac{m+1}{n} \pi\right) \sin \left(\frac{\pi}{n}\right)}{\sin ^{2}\left(\frac{\pi m}{2 n}\right) \sin ^{2}\left(\frac{(m+2)}{2 n} \pi\right)} \tag{4.7}
\end{equation*}
$$

Lemma 4.1 is now used with $a=2 n, b=m$ and $c=m+2$. This gives

$$
\begin{equation*}
c_{1}-b_{1}=\frac{\pi}{n} \text { and } c_{1}+b_{1}=\frac{\pi}{n}(m+1) \tag{4.8}
\end{equation*}
$$

to produce the result.
Example 4.3. Entry 4.235.3 states that

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\left(1-x^{2}\right) x^{n-3}}{1-x^{2 n}} \ln x d x=-\frac{\pi^{2}}{4 n^{2}} \tan ^{2} \frac{\pi}{n} \tag{4.9}
\end{equation*}
$$

The values $a=2 n, b=n-2$ and $c=n$ give

$$
\begin{equation*}
b_{1}=\frac{\pi}{2}-\frac{\pi}{n} \text { and } c_{1}=\frac{\pi}{2} \tag{4.10}
\end{equation*}
$$

This verifies the claim.
Example 4.4. Entry 4.235.4 appears as

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{m-1}+x^{n-m-1}}{1-x^{n}} \ln x d x=-\frac{\pi^{2}}{n^{2} \sin ^{2} \frac{\pi m}{n}} \tag{4.11}
\end{equation*}
$$

The change of variables $t=1 / x$ shows that the integral over $[1, \infty)$ is equal to that over $[0,1]$, therefore this entry should be written as

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{m-1}+x^{n-m-1}}{1-x^{n}} \ln x d x=-\frac{2 \pi^{2}}{n^{2} \sin ^{2} \frac{\pi m}{n}} \tag{4.12}
\end{equation*}
$$

to be consistent with the other entries in this section. The proof comes from Lemma 4.1 with $a=n, b=m$ and $c=n-m$.

## 5. Some singular integrals

The table [5] contains a variety of singular integrals of the form being discussed here. The examples considered in this section are evaluated employing the formula

$$
\begin{equation*}
\int_{0}^{\infty} \frac{t^{\mu-1} d t}{1-t}=\pi \cot \pi \mu \tag{5.1}
\end{equation*}
$$

To verify this evaluation, transform the integral over $[1, \infty)$ to $[0,1]$ by the change of variables $x \mapsto 1 / x$. This gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{t^{\mu-1} d t}{1-t}=\int_{0}^{1} \frac{t^{\mu-1}-t^{-\mu}}{1-t} d t \tag{5.2}
\end{equation*}
$$

This is entry $\mathbf{3 . 2 3 1}$.1. It was established in [7].
Differentiating with respect to $\mu$, the formula (5.1) gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{t^{\mu-1} \ln t}{1-t} d t=-\frac{\pi^{2}}{\sin ^{2} \pi \mu} \tag{5.3}
\end{equation*}
$$

and the change of variables $t=x^{a}$ gives

$$
\begin{equation*}
\omega(a, b):=\int_{0}^{\infty} \frac{x^{b-1} \ln x}{1-x^{a}} d x=-\frac{\pi^{2}}{a^{2} \sin ^{2}\left(\frac{\pi b}{a}\right)} . \tag{5.4}
\end{equation*}
$$

Example 5.1. Entry 4.251.2 states that

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{a-x}=\pi a^{\mu-1}\left(\ln a \cot (\pi \mu)-\frac{\pi}{\sin ^{2} \pi \mu}\right) \tag{5.5}
\end{equation*}
$$

The change of variables $x=a t$ yields

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{a-x}=a^{\mu-1} \int_{0}^{\infty} \frac{t^{\mu-1} \ln t}{1-t} d t+a^{\mu-1} \ln a \int_{0}^{\infty} \frac{t^{\mu-1} d t}{1-t} \tag{5.6}
\end{equation*}
$$

The result now follows from (5.1) and (5.3). It is probably clearer to write this entry as

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{a-x}=\pi a^{\mu-1}\left(\frac{\ln a}{\tan \pi \mu}-\frac{\pi}{\sin ^{2} \pi \mu}\right) \tag{5.7}
\end{equation*}
$$

to avoid possible confusions.

Example 5.2. Entry 4.252 .3 is

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1} \ln x}{1-x^{2}} d x=-\frac{\pi^{2}}{4} \operatorname{cosec}^{2} \frac{\pi p}{2} \tag{5.8}
\end{equation*}
$$

This is $\omega(2, p)$ and the result follows from (5.4).
Example 5.3. Entry 4.255.3 states that

$$
\begin{equation*}
\int_{0}^{\infty} \frac{1-x^{p}}{1-x^{2}} \ln x d x=\frac{\pi^{2}}{4} \tan ^{2}\left(\frac{\pi p}{2}\right) \tag{5.9}
\end{equation*}
$$

This is $\omega(1,2)-\omega(p+1,2)$ and the result comes from (5.4).
Example 5.4. Entry 4.252 .1 is written as

$$
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x d x}{(x+a)(x+b)}=\frac{\pi}{(b-a) \sin \pi \mu}\left[a^{\mu-1} \ln a-b^{\mu-1} \ln b-\pi \frac{a^{\mu-1}-b^{\mu-1}}{\tan \pi \mu}\right]
$$

This value follows from the partial fraction decomposition

$$
\begin{equation*}
\frac{1}{(x+a)(x+b)}=\frac{1}{b-a} \frac{1}{x+a}-\frac{1}{b-a} \frac{1}{x+b} \tag{5.10}
\end{equation*}
$$

and entry 4.251.1

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{x+c} d x=\frac{\pi c^{\mu-1}}{\sin \pi \mu}(\ln c-\pi \cot \pi \mu) \tag{5.11}
\end{equation*}
$$

established in [11]. Differentiating (5.11) with respect to $c$ yields

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{(x+c)^{2}} d x=-\frac{(\mu-1) c^{\mu-2} \pi}{\sin \pi \mu}\left(\ln c-\pi \cot \pi \mu+\frac{1}{\mu-1}\right) \tag{5.12}
\end{equation*}
$$

This is entry 4.252 .4 .
Example 5.5. Entry 4.257.1

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu} \ln (x / a) d x}{(x+a)(x+b)}=\frac{\pi\left[b^{\mu} \ln (b / a)+\pi\left(a^{\mu}-b^{\mu}\right) \cot \pi \mu\right]}{(b-a) \sin \pi \mu} \tag{5.13}
\end{equation*}
$$

follows from (5.11) and the beta integral

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{\mu-1} d x}{x+a}=\frac{\pi a^{\mu-1}}{\sin \pi \mu} \tag{5.14}
\end{equation*}
$$

This appears as entry $\mathbf{3 . 1 9 4 . 3}$ and it was established in [11].
Example 5.6. The change of variables $t=x^{q}$ gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1} d x}{1-x^{q}}=\frac{1}{q} \int_{0}^{\infty} \frac{t^{p / q-1} d x}{1-t}=\frac{\pi}{q} \cot \left(\frac{\pi p}{q}\right) \tag{5.15}
\end{equation*}
$$

from (5.3). This is entry $\mathbf{3 . 2 4 1 . 3}$. The special case $q=1$ gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1} d x}{1-x}=\pi \cot \pi p \tag{5.16}
\end{equation*}
$$

Differentiating with respect to $p$ produces

$$
\begin{equation*}
\int_{0}^{\infty} \frac{x^{p-1} \ln x}{1-x} d x=-\frac{\pi^{2}}{\sin ^{2} \pi p} \tag{5.17}
\end{equation*}
$$

The partial fraction decomposition

$$
\begin{equation*}
\frac{1}{(x+a)(x-1)}=\frac{1}{a+1} \frac{1}{x-1}-\frac{1}{a+1} \frac{1}{x+a} \tag{5.18}
\end{equation*}
$$

then produces entry $\mathbf{4 . 2 5 2 . 2}$
(5.19) $\int_{0}^{\infty} \frac{x^{\mu-1} \ln x}{(x+a)(x-1)} d x=\frac{\pi}{(a+1) \sin ^{2} \pi \mu}\left[\pi-a^{\mu-1}(\ln a \sin \pi \mu-\pi \cos \pi \mu)\right]$.

Example 5.7. The change of variables $t=x^{q}$ produces

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\ln x d x}{x^{p}\left(x^{q}-1\right)}=-\frac{1}{q^{2}} \int_{0}^{\infty} \frac{t^{(1-p) / q-1} \ln t d t}{1-t} \tag{5.20}
\end{equation*}
$$

Then, (5.3) gives

$$
\begin{equation*}
\int_{0}^{\infty} \frac{\ln x d x}{x^{p}\left(x^{q}-1\right)}=\frac{\pi^{2}}{q^{2}} \frac{1}{\sin ^{2}\left(\frac{p-1}{q} \pi\right)} \tag{5.21}
\end{equation*}
$$

This is entry 4.254 .3 .
Example 5.8. Entry 4.255.2 is

$$
\begin{equation*}
\int_{0}^{1} \frac{\left(1+x^{2}\right) x^{p-2}}{1-x^{2 p}} \ln x d x=-\left(\frac{\pi}{2 p}\right)^{2} \sec ^{2} \frac{\pi}{2 p} \tag{5.22}
\end{equation*}
$$

The evaluation of this entry starts with entry $\mathbf{3 . 2 3 1 . 5}$

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{\mu-1}-x^{\nu-1}}{1-x} d x=-\psi(\mu)+\psi(\nu) \tag{5.23}
\end{equation*}
$$

that was establsihed in [7]. The special case $\mu=1$

$$
\begin{equation*}
\int_{0}^{1} \frac{1-x^{\nu-1}}{1-x} d x=-\psi(1)+\psi(\nu) \tag{5.24}
\end{equation*}
$$

is differentiated with respect to $\nu$ to produce

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{\nu-1} \ln x}{1-x} d x=-\psi^{\prime}(\nu) \tag{5.25}
\end{equation*}
$$

The change of variables $x=t^{b}$ gives

$$
\begin{equation*}
\int_{0}^{1} \frac{t^{c-1} \ln t}{1-t^{b}} d t=-\frac{1}{b^{2}} \psi^{\prime}\left(\frac{c}{b}\right) \tag{5.26}
\end{equation*}
$$

Therefore

$$
\begin{aligned}
\int_{0}^{1} \frac{\left(1-x^{2}\right) x^{p-2}}{1-x^{2 p}} \ln x d x & =\int_{0}^{1} \frac{x^{p-2}}{1-x^{2 p}} \ln x d x+\int_{0}^{1} \frac{x^{p}}{1-x^{2 p}} \ln x d x \\
& =-\frac{1}{4 p^{2}}\left[\psi^{\prime}\left(\frac{1}{2}-\frac{1}{2 p}\right)+\psi^{\prime}\left(\frac{1}{2}+\frac{1}{2 p}\right)\right]
\end{aligned}
$$

The result now follows from the reflection formula for the polygamma function $\psi^{\prime}$ given in (2.14).

## 6. Combinations of logarithms and algebraic functions

This section presents the evaluation of some entries in [5] of the form

$$
\begin{equation*}
\int_{a}^{b} E_{1}(x) \ln E_{2}(x) d x \tag{6.1}
\end{equation*}
$$

where $E_{1}$ or $E_{2}$ is an algebraic function. Some of these have appeared in previous papers in this series. For example, entry 4.241.11

$$
\begin{equation*}
\int_{0}^{1} \frac{\ln x d x}{\sqrt{x\left(1-x^{2}\right)}}=-\frac{\sqrt{2 \pi}}{8} \Gamma^{2}\left(\frac{1}{4}\right) \tag{6.2}
\end{equation*}
$$

and entry 4.241 .5

$$
\begin{equation*}
\int_{0}^{1} \ln x \sqrt{\left(1-x^{2}\right)^{2 n-1}} d x=-\frac{(2 n-1)!!}{4(2 n)!!} \pi[\psi(n+1)+\gamma+\ln 4] \tag{6.3}
\end{equation*}
$$

were evaluated in [7]. Here $\psi(x)$ is the digamma function and $\gamma$ is Euler's constant.
Note 6.1. Define the family of integrals

$$
\begin{equation*}
f_{n}(a):=\int_{0}^{1} \frac{x^{a} \ln ^{n} x d x}{\sqrt{1-x^{2}}} \tag{6.4}
\end{equation*}
$$

Special cases include entry 4.241.7

$$
\begin{equation*}
\int_{0}^{1} \frac{\ln x d x}{\sqrt{1-x^{2}}}=-\frac{\pi}{2} \ln 2 \tag{6.5}
\end{equation*}
$$

that was evaluated in $[7]$ and entry $\mathbf{4 . 2 6 1 . 9}$

$$
\begin{equation*}
\int_{0}^{1} \frac{\ln ^{2} x d x}{\sqrt{1-x^{2}}}=\frac{\pi}{2}\left(\ln ^{2} 2+\frac{\pi^{2}}{12}\right) \tag{6.6}
\end{equation*}
$$

A trigonometric form of the family is obtained by the change of variables $x=\sin t$ :

$$
\begin{equation*}
f_{n}(a)=\int_{0}^{\pi / 2} \sin ^{a} t \ln ^{n} \sin t d t \tag{6.7}
\end{equation*}
$$

Theorem 6.2. The integral $f_{n}(a)$ is given by

$$
\begin{equation*}
f_{n}(a)=\lim _{s \rightarrow a}\left(\frac{d}{d s}\right)^{n} h(s) \tag{6.8}
\end{equation*}
$$

where

$$
\begin{equation*}
h(s)=\int_{0}^{\pi / 2} \sin ^{s} t d t=\frac{1}{2} B\left(\frac{s+1}{2}, \frac{1}{2}\right)=\frac{\sqrt{\pi}}{2} \frac{\Gamma\left(\frac{s}{2}+\frac{1}{2}\right)}{\Gamma\left(\frac{s}{2}+1\right)} . \tag{6.9}
\end{equation*}
$$

This appears as entry $\mathbf{3 . 6 2 1 . 5}$. Therefore, the evaluation of $f_{n}(a)$ requires the values of $\Gamma^{(k)}(x)$ for $0 \leqslant k \leqslant n$ at $x=(a+1) / 2$ and $x=a / 2+1$.

Example 6.3. For example,

$$
\begin{aligned}
f_{1}(0)=\int_{0}^{1} \frac{\ln x d x}{\sqrt{1-x^{2}}} & =\lim _{s \rightarrow 0} \frac{d}{d s}\left[\frac{\sqrt{\pi}}{2} \frac{\Gamma\left(\frac{s}{2}+\frac{1}{2}\right)}{\Gamma\left(\frac{s}{2}+1\right)}\right] \\
& =\frac{\sqrt{\pi}}{4} \frac{\Gamma^{\prime}(1 / 2) \Gamma(1)-\Gamma^{\prime}(1) \Gamma(1 / 2)}{\Gamma^{2}(1)}
\end{aligned}
$$

The values

$$
\begin{equation*}
\Gamma^{\prime}\left(\frac{1}{2}\right)=-\sqrt{\pi}(\gamma+2 \ln 2), \Gamma^{\prime}(1)=-\gamma, \Gamma\left(\frac{1}{2}\right)=\sqrt{\pi} \text { and } \Gamma(1)=1 \tag{6.10}
\end{equation*}
$$

give

$$
\begin{equation*}
f_{1}(0)=-\frac{\pi}{2} \ln 2 \tag{6.11}
\end{equation*}
$$

Proposition 6.4. The derivatives of the gamma function satisfy the recurrence

$$
\begin{equation*}
\Gamma^{(n+1)}(x)=\sum_{k=0}^{n}\binom{n}{k} \Gamma^{(k)}(x) \psi^{(n-k)}(x) . \tag{6.12}
\end{equation*}
$$

Example 6.5. A direct application of formula (6.8) evaluates entry 4.261.9

$$
\begin{equation*}
f_{2}(0)=\int_{0}^{1} \frac{\ln ^{2} x d x}{\sqrt{1-x^{2}}} \tag{6.13}
\end{equation*}
$$

Indeed, using $\Gamma(1)=1$, gives

$$
\begin{equation*}
f_{2}(0)=\frac{\sqrt{\pi}}{2}\left[-\frac{1}{2} \Gamma^{\prime}\left(\frac{1}{2}\right) \Gamma^{\prime}(1)+\frac{1}{2} \Gamma\left(\frac{1}{2}\right) \Gamma^{\prime}(1)^{2}+\frac{1}{4} \Gamma^{\prime \prime}\left(\frac{1}{2}\right)-\frac{1}{4} \Gamma\left(\frac{1}{2}\right) \Gamma^{\prime \prime}(1)\right] \tag{6.14}
\end{equation*}
$$

The values

$$
\begin{equation*}
\Gamma^{\prime \prime}(1)=\gamma^{2}+\frac{\pi^{2}}{6} \text { and } \Gamma^{\prime \prime}\left(\frac{1}{2}\right)=\frac{1}{2} \pi^{5 / 2}+\sqrt{\pi}(\gamma+2 \ln 2)^{2} \tag{6.15}
\end{equation*}
$$

give the identity (6.6).
It remains to explain the values given in (6.10) and (6.15). The recurrence (6.12) reduces the computation of the derivatives of $\Gamma(x)$ to those of $\psi(x)$. The special values given above come from the next result.

Lemma 6.6. The digamma function satisfies

$$
\begin{aligned}
\psi^{(n)}(1) & =(-1)^{n+1} n!\zeta(n+1) \\
\psi^{(n)}\left(\frac{1}{2}\right) & =(-1)^{n+1} n!\left(2^{n+1}-1\right) \zeta(n+1)
\end{aligned}
$$

Proof. This comes directly from (2.9).
Example 6.7. The values given in Lemma 6.6 yield

$$
\begin{aligned}
& f_{3}(0)=\int_{0}^{1} \frac{\ln ^{3} x d x}{\sqrt{1-x^{2}}}=-\frac{\pi}{8}\left(\pi^{2} \ln 2+4 \ln ^{3} 2+6 \zeta(3)\right) \\
& f_{4}(0)=\int_{0}^{1} \frac{\ln ^{4} x d x}{\sqrt{1-x^{2}}}=\frac{\pi}{480}\left(19 \pi^{4}+120 \pi^{2} \ln ^{2} 2+240 \ln ^{4} 2+1440 \ln 2 \zeta(3)\right)
\end{aligned}
$$

and

$$
\begin{aligned}
& f_{1}\left(\frac{1}{2}\right)=\int_{0}^{1} \frac{\sqrt{x} \ln x d x}{\sqrt{1-x^{2}}} \\
&=\frac{(\pi-4)}{\sqrt{2 \pi}} \Gamma^{2}\left(\frac{3}{4}\right) \\
& f_{2}\left(\frac{1}{2}\right)=\int_{0}^{1} \frac{\sqrt{x} \ln ^{2} x d x}{\sqrt{1-x^{2}}}=\frac{1}{2 \sqrt{2 \pi}} \Gamma^{2}\left(\frac{3}{4}\right)(32-16 G+\pi(\pi-8))
\end{aligned}
$$

where $G$ is Catalan's constant

$$
\begin{equation*}
G=\sum_{n=0}^{\infty} \frac{(-1)^{n}}{(2 n+1)^{2}} \tag{6.16}
\end{equation*}
$$

Example 6.8. Entry 4.261 .15 states that

$$
\begin{align*}
& \int_{0}^{1} \frac{x^{2 n} \ln ^{2} x}{\sqrt{1-x^{2}}} d x=  \tag{6.17}\\
& \qquad \frac{(2 n-1)!!}{2(2 n)!!} \pi\left\{\frac{\pi^{2}}{12}+\sum_{k=1}^{2 n} \frac{(-1)^{k}}{k^{2}}+\left[\sum_{k=1}^{2 n} \frac{(-1)^{k}}{k}+\ln 2\right]^{2}\right\} .
\end{align*}
$$

This is obtained by differentiating $h(s)$ twice with respect to $s$ to produce

$$
\begin{aligned}
& \int_{0}^{1} \frac{x^{s} \ln ^{2} x d x}{\sqrt{1-x^{2}}}= \\
& \quad \frac{\sqrt{\pi}}{8} \frac{\Gamma\left(\frac{s+1}{2}\right)}{\Gamma\left(\frac{s}{2}+1\right)}\left[\left(\psi\left(\frac{s}{2}+1\right)-\psi\left(\frac{s+1}{2}\right)\right)^{2}+\psi^{\prime}\left(\frac{s+1}{2}\right)-\psi^{\prime}\left(\frac{s}{2}+1\right)\right] .
\end{aligned}
$$

Therefore

$$
\int_{0}^{1} \frac{x^{2 n} \ln ^{2} x d x}{\sqrt{1-x^{2}}}=\frac{\sqrt{\pi}}{8} \frac{\Gamma\left(n+\frac{1}{2}\right)}{\Gamma(n+1)}\left[\left(\psi(n+1)-\psi\left(n+\frac{1}{2}\right)\right)^{2}+\psi^{\prime}\left(n+\frac{1}{2}\right)-\psi^{\prime}(n+1)\right] .
$$

The special values

$$
\begin{equation*}
\Gamma\left(n+\frac{1}{2}\right)=\frac{(2 n-1)!!}{2^{n}} \sqrt{\pi} \text { and } \Gamma(n+1)=n! \tag{6.18}
\end{equation*}
$$

give

$$
\int_{0}^{1} \frac{x^{2 n} \ln ^{2} x d x}{\sqrt{1-x^{2}}}=\frac{\pi}{8} \frac{(2 n-1)!!}{(2 n)!!}\left[\left(\psi(n+1)-\psi\left(n+\frac{1}{2}\right)^{2}+\psi^{\prime}\left(n+\frac{1}{2}\right)-\psi^{\prime}(n+1)\right] .\right.
$$

Now use the special values

$$
\begin{equation*}
\psi(n+1)=-\gamma+\sum_{k=1}^{n} \frac{1}{k} \text { and } \psi\left(n+\frac{1}{2}\right)=-\gamma-2 \ln 2+2 \sum_{k=1}^{n} \frac{1}{2 k-1} \tag{6.19}
\end{equation*}
$$

as well as

$$
\begin{equation*}
\psi^{\prime}(n+1)=\frac{\pi^{2}}{6}-\sum_{k=1}^{n} \frac{1}{k^{2}} \text { and } \psi^{\prime}\left(n+\frac{1}{2}\right)=\frac{\pi^{2}}{2}-4 \sum_{k=1}^{n} \frac{1}{(2 k-1)^{2}} \tag{6.20}
\end{equation*}
$$

to obtain

$$
\begin{equation*}
\psi(n+1)-\psi\left(n+\frac{1}{2}\right)=2 \sum_{k=1}^{2 n} \frac{(-1)^{k}}{k}+2 \ln 2 \tag{6.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\psi^{\prime}\left(n+\frac{1}{2}\right)-\psi^{\prime}(n+1)=\frac{\pi^{2}}{3}+4 \sum_{k=1}^{2 n} \frac{(-1)^{k}}{k^{2}} . \tag{6.22}
\end{equation*}
$$

This gives the result.
Example 6.9. A similar analysis gives entry 4.261.16

$$
\begin{aligned}
\int_{0}^{1} \frac{x^{2 n+1} \ln ^{2} x}{\sqrt{1-x^{2}}} d x= & \\
& -\frac{(2 n)!!}{(2 n+1)!!}\left\{\frac{\pi^{2}}{12}+\sum_{k=1}^{2 n+1} \frac{(-1)^{k}}{k^{2}}-\left[\sum_{k=1}^{2 n+1} \frac{(-1)^{k}}{k}+\ln 2\right]^{2}\right\}
\end{aligned}
$$

Example 6.10. Entry 4.241.6 states that

$$
\begin{equation*}
\int_{0}^{1 / \sqrt{2}} \frac{\ln x d x}{\sqrt{1-x^{2}}}=-\frac{\pi}{4} \ln 2-\frac{G}{2} \tag{6.23}
\end{equation*}
$$

The change of variables $x=\sin t$ gives

$$
\begin{equation*}
\int_{0}^{1 / \sqrt{2}} \frac{\ln x d x}{\sqrt{1-x^{2}}}=\int_{0}^{\pi / 4} \ln \sin t d t \tag{6.24}
\end{equation*}
$$

This integral is entry 4.224 .2 and it has been evaluated in [3].

## 7. An example producing a trigonometric answer

The next example contains, in the logarithmic part, a quotient of linear functions. The evaluation of this entry requires a different approach.

Example 7.1. Entry 4.297 .8 states that

$$
\begin{equation*}
\int_{0}^{1} \ln \frac{1+a x}{1-a x} \frac{d x}{x \sqrt{1-x^{2}}}=\pi \sin ^{-1} a . \tag{7.1}
\end{equation*}
$$

This evaluation starts with the expansion

$$
\begin{equation*}
\frac{1}{x} \ln \frac{1+a x}{1-a x}=\sum_{n=0}^{\infty} \frac{2 a^{2 n+1}}{2 n+1} x^{2 n} \tag{7.2}
\end{equation*}
$$

to obtain

$$
\begin{equation*}
\int_{0}^{1} \ln \frac{1+a x}{1-a x} \frac{d x}{x \sqrt{1-x^{2}}}=\sum_{n=0}^{\infty} \frac{2 a^{2 n+1}}{2 n+1} \int_{0}^{1} \frac{x^{2 n} d x}{\sqrt{1-x^{2}}} \tag{7.3}
\end{equation*}
$$

The change of variables $x=\sin \theta$ gives

$$
\begin{equation*}
\int_{0}^{1} \frac{x^{2 n} d x}{\sqrt{1-x^{2}}}=\int_{0}^{\pi / 2} \sin ^{2 n} \theta d \theta=\frac{\pi}{2^{2 n+1}}\binom{2 n}{n} \tag{7.4}
\end{equation*}
$$

The last evaluation is the famous Wallis' formula. It appears as entry $\mathbf{3 . 6 2 1 . 3}$ and it was established in [2] and [12]. Therefore

$$
\begin{equation*}
\int_{0}^{1} \ln \frac{1+a x}{1-a x} \frac{d x}{x \sqrt{1-x^{2}}}=\sum_{n=0}^{\infty} \frac{\pi}{2^{2 n}} \frac{a^{2 n+1}}{2 n+1}\binom{2 n}{n} . \tag{7.5}
\end{equation*}
$$

The series is now identified from the classical expansion

$$
\begin{aligned}
\sin ^{-1} x & =\sum_{n=0}^{\infty} \frac{\left(\frac{1}{2}\right)_{n}}{(2 n+1) n!} x^{2 n+1} \\
& =\sum_{n=0}^{\infty} \frac{1}{2^{2 n}(2 n+1)}\binom{2 n}{n} x^{2 n+1}
\end{aligned}
$$

obtained by expanding the integrand in

$$
\begin{equation*}
\sin ^{-1} x=\int_{0}^{x} \frac{d t}{\sqrt{1-t^{2}}} \tag{7.6}
\end{equation*}
$$

as a binomial series and integrating term by term.

Further examples in [5], of the class considered here, will be presented in a future publication.

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