THE INTEGRALS IN GRADSHTEYN AND RHYZIK. PART 3: COMBINATIONS OF LOGARITHMS AND EXPONENTIALS.

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ABSTRACT. We present the evaluation of a family of exponential-logarithmic integrals. These have integrands of the form $P(e^{tx}, \ln x)$ where P is a polynomial. The examples presented here appear in sections 4.33, 4.34 and 4.35 in the classical table of integrals by I. Gradshteyn and I. Rhyzik.

1. Introduction

This is the third in a series of papers dealing with the evaluation of definite integrals in the table of Gradshteyn and Rhyzik [2]. We consider here problems of the form

(1.1)
$$\int_0^\infty e^{-tx} P(\ln x) \, dx,$$

where t>0 is a parameter and P is a polynomial. In future work we deal with the finite interval case

(1.2)
$$\int_a^b e^{-tx} P(\ln x) dx,$$

where $a, b \in \mathbb{R}^+$ with a < b and $t \in \mathbb{R}$. The classical example

(1.3)
$$\int_0^\infty e^{-x} \ln x \, dx = -\gamma,$$

where γ is Euler's constant is part of this family. The integrals of type (1.1) are linear combinations of

(1.4)
$$J_n(t) := \int_0^\infty e^{-tx} (\ln x)^n dx.$$

The values of these integrals are expressed in terms of the gamma function

(1.5)
$$\Gamma(s) = \int_0^\infty x^{s-1} e^{-x} dx$$

and its derivatives.

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2. The evaluation

In this section we consider the value of $J_n(t)$ defined in (1.4). The change of variables s = tx yields

(2.1)
$$J_n(t) = \frac{1}{t} \int_0^\infty e^{-s} (\ln s - \ln t)^n ds.$$

Expanding the power yields J_n as a linear combination of

(2.2)
$$I_n := \int_0^\infty e^{-x} (\ln x)^n \, dx.$$

An analytic expression for these integrals can be obtained directly from the representation of the $gamma\ function$ in (1.5).

Proposition 2.1. For $n \in \mathbb{N}$ we have

(2.3)
$$\int_0^\infty (\ln x)^n \ x^{s-1} e^{-x} \, dx = \left(\frac{d}{ds}\right)^n \Gamma(s).$$

In particular

(2.4)
$$I_n := \int_0^\infty (\ln x)^n e^{-x} dx = \Gamma^{(n)}(1).$$

Proof. Differentiate (1.5) n-times with respect to the parameter s.

Example 2.2. Formula 4.331.1 in [2] states that¹

(2.5)
$$\int_0^\infty e^{-\mu x} \ln x \, dx = -\frac{\delta}{\mu}.$$

where $\delta = \gamma + \ln \mu$. This value follows directly by the change of variables $s = \mu x$ and the classical special value $\Gamma'(1) = -\gamma$. The reader will find in chapter 9 of [1] details on this constant. In particular, if $\mu = 1$, then $\delta = \gamma$ and we obtain (1.3):

(2.6)
$$\int_0^\infty e^{-x} \ln x \, dx = -\gamma.$$

The change of variables $x = e^{-t}$ yields the form

(2.7)
$$\int_{-\infty}^{\infty} t e^{-t} e^{-e^{-t}} dt = \gamma.$$

Many of the evaluations are given in terms of the polygamma function

(2.8)
$$\psi(x) = \frac{d}{dx} \ln \Gamma(x).$$

Properties of ψ are summarized in Chapter 1 of [4]. A simple representation is

(2.9)
$$\psi(x) = \lim_{n \to \infty} \left(\ln n - \sum_{k=0}^{n} \frac{1}{x+k} \right),$$

from where we conclude that

(2.10)
$$\psi(1) = \lim_{n \to \infty} \left(\ln n - \sum_{k=0}^{n} \frac{1}{1+k} \right) = -\gamma,$$

¹The table uses C for the Euler constant.

this being the most common definition of the Euler's constant γ . This is precisely $\Gamma'(1) = -\gamma$.

The derivatives of ψ satisfy

(2.11)
$$\psi^{(m)}(x) = (-1)^{m+1} m! \, \zeta(m+1, x),$$

where

(2.12)
$$\zeta(z,q) := \sum_{n=0}^{\infty} \frac{1}{(n+q)^z}$$

is the *Hurwitz zeta function*. This function appeared in [3] in the evaluation of some logarithmic integrals.

Example 2.3. Formula 4.335.1 in [2] states that

(2.13)
$$\int_0^\infty e^{-\mu x} (\ln x)^2 dx = \frac{1}{\mu} \left[\frac{\pi^2}{6} + \delta^2 \right].$$

This can be verified using the procedure described above: the change of variable $s = \mu x$ yields

(2.14)
$$\int_0^\infty e^{-\mu x} (\ln x)^2 dx = \frac{1}{\mu} (I_2 - 2 \ln \mu I_1 + \ln^2 \mu I_0),$$

where I_n is defined in (2.4). To complete the evaluation we need some special values: $\Gamma(1) = 1$ is elementary, $\Gamma'(1) = \psi(1) = -\gamma$ appeared above and using (2.11) we have

(2.15)
$$\psi'(x) = \frac{\Gamma''(x)}{\Gamma(x)} - \left(\frac{\Gamma'(x)}{\Gamma(x)}\right)^2.$$

The value

(2.16)
$$\psi'(1) = \zeta(2) = \frac{\pi^2}{6},$$

where $\zeta(z)=\zeta(z,1)$ is the Riemann zeta function. This comes directly from (2.11). Thus

(2.17)
$$\Gamma''(1) = \zeta(2) + \gamma^2.$$

The special case $\mu = 1$ produces

(2.18)
$$\int_0^\infty e^{-x} (\ln x)^2 dx = \zeta(2) + \gamma^2.$$

Similar arguments yields formula **4.335.3** in [2]:

(2.19)
$$\int_0^\infty e^{-\mu x} (\ln x)^3 dx = -\frac{1}{\mu} \left[\delta^3 + \frac{1}{2} \pi^2 \delta - \psi''(1) \right],$$

with the special case $\mu = 1$ now yielding

(2.20)
$$\int_0^\infty e^{-x} (\ln x)^3 dx = -\gamma^3 - \frac{1}{2} \pi^2 \gamma + \psi''(1).$$

Problem 2.4. In [1], page 203, we introduced the notion of *weight* for some real numbers. In particular, we have assigned $\zeta(j)$ the weight j. Differentiation increases the weight by 1, so that $\zeta'(3)$ has weight 4. The task is to check that the integral

(2.21)
$$I_n := \int_0^\infty e^{-x} (\ln x)^n dx$$

is a homogeneous form of weight n.

3. A SMALL VARIATION

Similar arguments are now employed to produce a larger family of integrals. The representation

(3.1)
$$\int_{0}^{\infty} x^{s-1} e^{-\mu x} dx = \mu^{-s} \Gamma(s),$$

is differentiated n times with respect to the parameter s to produce

(3.2)
$$\int_0^\infty (\ln x)^n x^{s-1} e^{-\mu x} dx = \left(\frac{d}{ds}\right)^n \left[\mu^{-s} \Gamma(s)\right].$$

The special case n = 1 yields

(3.3)
$$\int_0^\infty x^{s-1} e^{-\mu x} \ln x \, dx = \frac{d}{ds} \left[\mu^{-s} \Gamma(s) \right]$$
$$= \mu^{-s} \left(\Gamma'(s) - \ln \mu \Gamma(s) \right)$$
$$= \mu^{-s} \Gamma(s) \left(\psi(s) - \ln \mu \right).$$

This evaluation appears as **4.352.1** in [2]. The special case $\mu = 1$ yields

(3.4)
$$\int_0^\infty x^{s-1} e^{-x} \ln x \, dx = \Gamma'(s),$$

that is 4.352.4 in [2].

Special values of the gamma function and its derivatives yield more concrete evaluations. For example, the functional equation

(3.5)
$$\psi(x+1) = \psi(x) + \frac{1}{x},$$

that is a direct consequence of $\Gamma(x+1) = x\Gamma(x)$, yields

(3.6)
$$\psi(n+1) = -\gamma + \sum_{k=1}^{n} \frac{1}{k}.$$

Replacing in (3.3) we obtain

(3.7)
$$\int_0^\infty x^n e^{-\mu x} \ln x \, dx = \frac{n!}{\mu^{n+1}} \left(\sum_{k=1}^n \frac{1}{k} - \gamma - \ln \mu \right),$$

that is 4.352.2 in [2].

The final formula of Section 4.352 in [2] is 4.352.3

$$\int_0^\infty x^{n-1/2} e^{-\mu x} \ln x \, dx = \frac{\sqrt{\pi} (2n-1)!!}{2^{\mu} \mu^{n+1/2}} \left[2 \sum_{k=1}^n \frac{1}{2k-1} - \gamma - \ln(4\mu) \right].$$

This can also be obtained from (3.3) by using the classical values

$$\Gamma(n + \frac{1}{2}) = \frac{\sqrt{\pi}}{2^n} (2n - 1)!!$$

$$\psi(n + \frac{1}{2}) = -\gamma + 2 \left(\sum_{k=1}^n \frac{1}{2k - 1} - \ln 2 \right).$$

The details are left to the reader.

Section 4.353 of [2] contains three peculiar combinations of integrands. The first two of them can be verified by the methods described above: formula 4.353.1 states

(3.8)
$$\int_0^\infty (x - \nu) x^{\nu - 1} e^{-x} \ln x \, dx = \Gamma(\nu),$$

and 4.353.2 is

(3.9)
$$\int_0^\infty (\mu x - n - \frac{1}{2}) x^{n - \frac{1}{2}} e^{-\mu x} \ln x \, dx = \frac{(2n - 1)!!}{(2\mu)^n} \sqrt{\frac{\pi}{\mu}}.$$

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